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# ON SYSTEMS OF ORTHOGONAL POLYNOMIALS WITH INNER AND END POINT MASSES

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Abstract. A system of orthogonal polynomials related to the sieved ultraspherical polynomials is presented. It is shown that their orthogonality measure carries end point masses. It is also proved that the orthogonality measure of their first associated (numerator) polynomials has an embedded masspoint.

Sumario. Se describe un sistema de polinomios ortogonales, relacionado con el de los ultraesféricos cribados, cuya medida de ortogonalidad tiene masas en los extremos del intervalo de ortogonalidad. Se demuestra también que la medida de ortogonalidad de los polinomios numeradores correspondientes presenta una masa sumergida.

**§1.** Introduction. We study in this paper two systems of orthogonal polynomials defined by recurrence relations of the form

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Key words and Phrases. Orthogonal polynomials, recurrence relations, continued fractions, orthogonality measures, mass points.

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$$xP_{2n}(x) = P_{2n+1}(x) + a_{0,n}P_{2n-1}(x)$$

$$xP_{2n+1}(x) = P_{2n+2}(x) + a_{1,n}P_{2n}(x), \quad n \ge 0,$$

and initial conditions

$$P_{-1}(x) = 0, P_{\hat{0}}(x) = 1.$$
 (1.2)

We take

$$a_{0,0} = a_{1,0} = \alpha/2$$
 (1.3)  
 $a_{0,n} = n/(4(n+\tau)), a_{1,n} = (n+2\tau)/(4(n+\tau)), n \ge 1$ 

and

$$a_{0,0} = \alpha/2, a_{1,0} = 1/(4(1+\tau))$$
 (1.4)

$$a_{0,n} = (n+2\tau)/(4(n+\tau)), a_{1,n} = (n+1)/(4(n+\tau+1)), n \ge 1.$$

In both cases  $\alpha, \tau > 0$ . We observe that the polynomials determined by (1.4) are the first associated of those defined by (1.3). Hence, we will denote with  $\{P_n(x)\}$  the polynomials given by (1.3) and with  $\{P_n^{\{1\}}(x)\}$  those by (1.4).

Recurrence relations of the form (1.1) have frequently appeared in the literature of the last few years (Chihara [9], Charris and Gómez [6]) and are related to spectral problems in quantum chemistry (Slim [23], Wheeler [27]). They are also special cases of sieved orthogonal polynomials (A1-Salam, Allaway and Askey [2], Charris and Ismail [7], [8], Ismail [12], [13], [14]). As a matter of fact, the system  $\{P_n(x)\}$  described above reduces to sieved ultraspherical polynomials of the first kind when  $\alpha = 1$  (see [2]), and  $\{P_n^{(1)}(x)\}$  is then the system of their first associated (observe that  $\{P_n^{(1)}(x)\}$  is independent of  $\alpha$ ). In [2] it is shown that, when  $\alpha = 1$ ,  $\{P_n(x)\}$  is a system of orthogonal polynomials with respect to the absolutely continuous measure

$$d\mu(x) = |x|^{2\tau} (1-x^2)^{\tau-\frac{1}{2}} \Phi(x) dx, \qquad (1.5)$$

where  $\Phi(x)$  is the characteristic function of the interval (-1,1), but the case  $\alpha \not= 1$  was not considered. Neither was considered the system  $\{P_n^{(1)}(x)\}$  of the first associated polynomials (the sieved ultraspherical polynomials of the second kind of Al-Salam, Allaway and Askey [2] are first associated polynomials of a system of orthogonal polynomials but not of the sieved ultraspherical of the first kind). The point of view adopted in [2] did not allow for the treatment of these cases. The purpose of the present paper is to establish some properties of  $\{P_n(x)\}$  and  $\{P_n^{(1)}(x)\}$  that seem to have gone unnoticed in the literature, perhaps because of the above reason.

First,  $\{P_n(x)\}$  has, for special values of a and  $\tau$ , masses at the end points of the interval of orthogonality. This phenomenum has been explored by Chihara [10], Koornwinder [18] and Koekoek [15], [16], [17], but their examples of this situation are rather complicated and artificial. Second, and even more remarkable, is the fact that  $\{P_n^{(1)}(x)\}$  may have masses which are embedded in the continuous spectrum, i.e., which are interior to the support of the orthogonality measure. Examples of this situation were, as far as we know, absent from the literature prior to Geronimo and Van Assche [11]. However, the examples in [11] are more sophisticated and not so explicit.

In Section 2 we review the basic notions about orthogonal polynomials and in Section 3 we give a brief account of ultraspherical polynomials and their associated. Section 4 contains some basic formulae, and Section 5 the main results. Finally, in Section 6 we study the connection with spectral theory and examine futher generalizations.

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bian Symposium in functional Analysis, Medellin, July 1989. It originated in the research of this author for his Masters thesis at the National University of Colombia in Bogotá.

- **52.** Background. A system  $\{P_n(x)/n \ge 0\}$  of polynomials with real coefficients is called a system of orthogonal poynomials (0.P.S.) if
  - 1)  $P_n(x)$  has degree n and  $P_0(x) = 1$ .
  - 2) There is a positive measure  $\mu$  supported by the real line such that

$$\int_{-\infty}^{+\infty} P_{n}(x) P_{m}(x) d\mu = k_{n} \delta_{nm}; m, n \ge 0, \quad k_{n} > 0.$$
 (2.1)

It can be shown (See [6],[9]) that a system  $\{P_n(x)\}$  is an O.P.S if and only if there are real numbers  $A_n, B_n, C_n, n > 0$ , such that

$$A_n A_{n+1} C_n > 0, \quad n \ge 0,$$
 (2.2)

$$P_{n+1}(x) = (A_n x + B_n) P_n(x) - C_n P_{n-1}(x), \quad n \ge 0,$$
 (2.3)

and

$$P_{-1}(x) = 0, \quad P_{0}(x) = 1.$$
 (2.4)

If that is the case  $\mu$  can be so chosen that

$$k_0 = 1; k_n = (A_0/A_n) C_1...C_n, \quad n \ge 1.$$
 (2.5)

The numbers  $A_n, B_n, C_n$  are uniquely determined by  $\{P_n(x)\}$ , except for  $C_0$ , which is arbitrary (Favard's Theorem).

The measure  $\mu$  satisfying (2.1) has bounded support, Supp $\mu$ , if and only if there is a constant M>0 such that

$$|B_n/A_n| \le M/3, \sqrt{C_{n+1}/A_nA_{n+1}} \le M/3, \quad n \ge 0,$$
 (2.6)

in which case Supp $\mu\subseteq \left[-M,M\right]$  and  $\mu$  is the only measure for

which (2.1) holds.

Let  $A_n, B_n, C_n$  be as in (2.3). The system  $\{P_n^{(i)}(x)\}, i = 0, 1, 2, 3, \ldots$ , defined by

$$P_{n+1}^{(i)}(x) = (A_{n+i}x + B_{n+i}) P_n^{(i)}(x) - C_{n+i}P_{n-1}^{(i)}(x)$$
 (2.7)

for n > 0 and

$$P_{-1}^{(i)}(x) = 0, \quad P_{0}^{(i)}(x) = 1$$
 (2.8)

is also orthogonal with respect to some positive measure  $\mu_{(i)}$ . It is called the system of  $i^{th}$ -associated polynomials of  $\{P_n(x)\}$ . Clearly  $P_n^{(0)}(x) = P_n(x)$ . In general  $\mu_{(i)} \neq \mu_{(j)}$  if  $i \neq j$ .

Let  $\{Q_n(x)\}$  be a system of polynomials with real coefficients such that (2.3) holds for  $\{Q_n(x)\}$  with  $n \ge 2$ . Then

$$Q_n(x) = (A - C/C_1)P_n(x) + (B + C/C_1P_1(x))P_{n-1}^{(1)}(x), n \ge 1, \qquad (2.9)$$

where

$$A = Q_{0}(x); B = Q_{1}(x) - Q_{0}(x)P_{1}(x),$$

$$C = Q_{2}(x) - P_{1}^{(1)}(x)Q_{1}(x) + C_{1}Q_{0}(x).$$
(2.10)

If (2.3) holds for  $\{Q_n(x)\}$  with  $n \ge 1$ , then C = 0; if it hols for  $n \ge 0$ , then B = C = 0; if in the latter case  $Q_0(x) = 1$ , then  $Q_n(x) = P_n(x)$  for all n.

Assume  $\{P_n(x)\}$  is an O.P.S. and let  $\mu$  be an orthogonality measure for  $\{P_n(x)\}$ . Let

$$p_n(x) = P_n(x) / \sqrt{k_n}, \quad n \geqslant 0;$$
 (2.11)

then

$$\int_{-\infty}^{+\infty} p_{n}(x) p_{m}(x) d\mu = \delta_{mn}; m, n > 0.$$
 (2.12)

The system  $\{p_n(x)\}$  is called the orthonormal polynomial system (O.N.P.S.) of  $\{P_n(x)\}$ . If  $\{P_n(x)\}$  satisfies (2.3) and (2.4) and  $k_n$  is given by (2.5), then  $\{p_n(x)\}$  satisfies the

recurrence relation

$$xp_n(x) = b_{n+1}p_{n+1}(x) + a_np_n(x) + b_np_{n-1}(x), \quad n \ge 1,$$
 (2.13)

and the initial conditions  $(x) = (x) \begin{pmatrix} x \\ x \end{pmatrix} \begin{pmatrix} x \\ x \end{pmatrix} = (x) \begin{pmatrix} x \\ x \end{pmatrix} \begin{pmatrix} x \\ x \end{pmatrix}$ 

$$p_0(x) = 1, p_1(x) = (1/b_1)(x-a_0),$$
 (2.14)

where

$$a_n = -B_n/A_n, b_{n+1} = \sqrt{C_{n+1}/A_nA_{n+1}}, \quad n \ge 0.$$
 (2.15)

Condition (2.6) translates into and believed all (1)4

$$|a_n| \le M/3, 0 < b_{n+1} \le M/3, \quad n \ge 0,$$
 (2.16)

in which case Supp $\mu \subseteq [-M,M]$ .

In what follows we will assume (2.6)(or, (2.16)) to hold. With  $\mu$  satisfying (2.1) we then have

$$\lim_{n \to \infty} A_0 \frac{P_{n-1}^{(1)}(x)}{P_n(x)} = \int_{-\infty}^{+\infty} \frac{d\mu(t)}{x - t}, \quad x \in \mathcal{C} - [-M, M]. \quad (2.17)$$

Relation (2.17) is known as Markov's Theorem. The left hand side of (2.17) is known as the continued fraction of  $\{P_n(x)\}$ . It is, actually, the limit of the continued fraction

$$A_0$$
 -  $C_1$  -  $C_2$  - ... (2.18)

The right hand side of (2.17) is  $X(x) = -2\pi i \hat{\mu}(x)$ , where  $\hat{\mu}(x)$ , called the Cauchy-Stieltjes transform of  $\mu$ , is an analytic function of x for  $x \not\in \text{Supp}\mu$ . The Stieltjes inversion formula

$$\int_{-\infty}^{+\infty} \int_{0}^{+\infty} \{(x) d\mu(x) = \lim_{\varepsilon \to 0+} \frac{1}{2\pi i} \int_{-\infty}^{+\infty} \{X(x - i\varepsilon) - X(x + i\varepsilon)\} \int_{0}^{+\infty} \{x dx, x \in \mathbb{R}, (2.19)\}$$

which holds for any continuous function f on  $\mathbb{R}$ , then allows to recover  $\mu$  from X(x). Let

$$\sigma(\tau) = \int_{-\infty}^{\tau} d\mu, \quad \tau \in \mathbb{R}.$$
 (2.20)

The right continuous, non-decreasing function  $\sigma$  is called the distribution function of  $\mu$ . From (2.19), we get

$$\sigma(\tau) = \lim_{\varepsilon \to 0^+} \frac{1}{2\pi i} \int_{-\infty}^{\tau} \{X(x - i\varepsilon) - X(x + i\varepsilon)\} dx, \quad \tau \in \mathbb{R}.$$
 (2.21)

If Supp $\mu \subseteq (a,b), -\infty < a < b < +\infty$ , then

$$\int_{a}^{b} f d\mu = \int_{a}^{b} f(x) d\sigma(x)$$
 (2.22)

where the right hand side is an ordinary Riemann-Stieltjes integral. Futhermore, Suppµ coincides with the set  $S_{\mu}$  of points of increase of  $\sigma$ , i.e., with the set of points  $x \in \mathbb{R}$  such that  $\sigma(x_1) < \sigma(x_2)$  if  $x_1 < x < x_2$ . We have  $S_{\mu} = C_{\mu} \cup P_{\mu}$ , where  $C_{\mu}$  is the set of points x in  $S_{\mu}$  such that  $\sigma(x) = \sigma(x-0)$ , and  $P_{\mu}$ , of those with  $\sigma(x) \neq \sigma(x-0)$  (this is a countable set). Clearly  $C_{\mu} \cap P_{\mu} = \phi$ . The set  $C_{\mu}$  is called the continuous support of  $\mu$ ; and  $P_{\mu}$ , the point support. We observe that they are not, in general, closed sets.

The following characterization of points in  $P_{\mu}$  is useful:  $x \in P_{\mu}$  if and only if

$$\sum_{n=0}^{\infty} \frac{p_n^2(x)}{k_n} < +\infty.$$
 (2.23)

In such case,

$$\mu(\{x\}) = \sigma(x) - \sigma(x-0) = 1 / \left(\sum_{n=0}^{\infty} \frac{p_n^2(x)}{k_n}\right). \tag{2.24}$$

Let  $\mathcal{D}_{\mu}$  be the set of isolated points in Supp $\mu$ . Clearly  $\mathcal{D}_{\mu}\subseteq\mathcal{P}_{\mu}$ . The points of  $S_{\mu}$  in  $\mathcal{D}_{\mu}$  are the isolated poles of X(x), which are all simple, and  $x\in\mathcal{D}_{\mu}$  if and only if

$$\mu(\{x\}) = \sigma(x) - \sigma(x - 0) = \text{Res}(X, x) \neq 0.$$
 (2.25)

Points in  $P_{\mu}$  are called mass points of  $\mu$ ; those in  $\mathcal{D}_{\mu}$ , isolated mass points. The points in  $P_{\mu}$  N interior  $(S_{\mu})$  are the embedded mass points. If  $S_{\mu}$  happens to be a union of intervals, the end points of the infinite intervals in this

class that are in  $P_{\mu}$  , but not in the interior of  $S_{\mu}$  , are the end-points masses of  $\mu.$ 

The measure µ can be written

$$\mu = \mu_c + \mu_s + \mu_f - 1(d, s) = uqque 11 (2.26)$$

where  $\mu_c$  is abolutely continuous,  $\mu_{\delta}$  is singular continuous and  $\mu_j$  is a jump measure. The measure  $\mu_c$  is concentrated on  $C_{\mu}$  and  $\mu_{\delta}$  on a subset of Lebesgue measure 0 of this set;  $\mu_j$  is concentrated on  $P_{\mu}$ . The measure  $\mu_c$  can be written  $g(x)\,dx$ , where g(x) is an integrable function (locally integrable if Supp $\mu$  is not assumed to be bounded).

**REMARK 2.1.** Let  $X_{\varepsilon}(x) = X(x-i\varepsilon) - X(x+i\varepsilon)$  and assume that

$$\lim_{\varepsilon \to 0} \frac{1}{2\pi i} X_{\varepsilon}(x)$$
 (2.27)

exists a.e. in an open interval (a,b). Further assume that for a < c < d < b there are C = C(c,d) > 0 and  $\varepsilon_0 = \varepsilon(c,d) > 0$  such that  $|X_{\varepsilon}(x)| \le C$  for almost all  $x \in [c,d]$  and all  $\varepsilon \le \varepsilon_0$ . Then  $g(x) = \lim_{\varepsilon \to 0} (1/2\pi \dot{\varepsilon}) \cdot X_{\varepsilon}(x)$ , a.e. on (a,b), and  $\mu$  is free of masses in (a,b). If  $g(x) \ne 0$  a.e. in (a,b) then  $(a,b) \subseteq C_{\mu}$ .

Proofs of Markov's theorem can be found in [1], [26]. See also [6]. The relevant facts about continued fractions and its relation to O.P.S. are in [3], [9], [26]. For the Cauchy-Sieltjes transform and the Stieltjes inversion formula see [1], [5], [6], [19]. A proof of relation (2.24) is in [1], [24], and one of (2.25), in [1], [6].

# §3. The Associated Ultraspherical Polynomials. The notation

$$(\alpha)_{n} = \begin{cases} 1, & n = 0 \\ \alpha(\alpha+1) \dots (\alpha+n-1), & n \geqslant 1, \end{cases}$$
 (3.1)

will be used in the sequel. Clearly

$$(\alpha)_n = \frac{\Gamma(\alpha+n)}{\Gamma(\alpha)}, \quad \alpha \neq 0,$$
 (3.2)

where  $\Gamma(x)$  denotes the Gamma function. The asymptotic formula

$$\frac{\Gamma(a+n)}{\Gamma(b+n)} \approx n^{a-b}, \quad n \to \infty$$
 (3.3)

follows from Stirling's formula ([20],[22]).

We also recall the binomial formula ([20]),[22])

$$(1-x)^{-\alpha} = \sum_{n=0}^{\infty} \frac{(\alpha)_n}{n!} \cdot x^n, \quad |x| < 1.$$
 (3.4)

The associated ultraspherical polynomials  $\{C_n^T(x;i)\}$  are given by

$$2(n+\tau+i)xC_n^{\tau}(x;i)$$

$$= (n+i+1)C_{n+1}^{\mathsf{T}}(x;i) + (n+2\tau+i-1)C_{n-1}^{\mathsf{T}}(x;i), \quad n \geq 0, \tag{3.5}$$

and the initial conditions

$$c_{-1}^{\tau}(x;i) = 0, \quad c_{0}^{\tau}(x;i) = 1.$$
 (3.6)

We assume  $i = 0, 1, 2, 3, \ldots$ , and write  $C_n^{\tau}(x) = C_n^{\tau}(x; 0)$ . The generating functions

$$\sum_{n=0}^{\infty} C_n^{\tau}(x) t^n = (t^2 - 2tx + 1)^{-\tau}$$
 (3.7)

and

$$\sum_{n=0}^{\infty} c_n^{\tau}(x;i) t^{n+1} = i(t^2 - 2tx + 1)^{-\tau} \int_0^t u^{i-1} (u^2 - 2ux + 1)^{\tau - 1} du, \quad i \ge 1, \quad (3.8)$$

are easily obtained from (3.5) and (3.6). Observe that

$$c_{n}^{\mathsf{T}}(x;i) = \sum_{k=0}^{n} \frac{i}{k+i} c_{k}^{\mathsf{1-T}}(x) c_{n-k}^{\mathsf{T}}(x), n \geqslant 0, i \geqslant 1, \tag{3.9}$$

follows from (3.8). Hence

$$C_n^{\tau}(1) = \frac{(2\tau)_n}{n!}, \quad n > 0,$$
 (3.10)

and also

$$c_{n-1}^{\tau}(1;1) = \frac{1}{2\tau - 1} \left[ \frac{(2\tau)_n}{n!} - 1 \right], \quad \tau \neq \frac{1}{2}$$

$$c_n^{\frac{1}{2}}(1;1) = \sum_{k=0}^{n} \frac{1}{k+1}, \quad n > 0.$$
(3.11)

The continued fraction of  $\{C_n^{\tau}(x)\}$  is

$$X_{o}(x) = 2\tau z \int_{0}^{1} (1 - \frac{z}{2} \cdot u)^{\tau - 1} (1 - u)^{\tau - 1} du$$
 (3.12)

where  $z=z(x)=x-\sqrt{x^2-1}$ ,  $z=z(x)=x+\sqrt{x^2-1}$  and  $\sqrt{x^2-1}$  is the branch of the square root of  $x^2-1$  which is analytic in  $\mathbb{C}$ -[-1,1] and behaves like x when  $x\to\infty$ . (It follows that  $|z(x)|\leqslant |z(x)|$  with |z(x)|=|z(x)| if and only if  $x\in[-1,1]$ , z=1/z, z+z=2x,  $z-z=2\sqrt{x^2-1}$ ). Formula (3.12) can be obtained from (3.7) and (3.8) by means of (3.4) and Darboux's asymptotic method ([21], Sec.8.9). Stieltjes inversion formula and Remark 2.1 then yield for the orthogonality measure  $\mu_0$  of  $\{C_n^{\tau}(x)\}$  the formula

$$d\mu_{o}(x) = -\frac{1}{\pi}\Phi(x) \operatorname{Im} X_{o}(x) dx$$

$$= \frac{2^{2\tau - 1}}{\pi} \frac{\Gamma(\tau)^{2}}{\Gamma(2\tau)} (1 - x^{2})^{\tau - \frac{1}{2}} \Phi(x) dx$$
(3.13)

where  $\phi(x)$  is the characteristic function of (-1,1). Hence,  $\mu_0$  is absolutely continuous.

The continued fraction of  $\{C_n^{\tau}(x;1)\}$  is

$$X_{1}(x) = 2(1+\tau)z \frac{\int_{0}^{1} u(1-\frac{z}{2}\cdot u)^{\tau-1}(1-u)^{\tau-1}du}{\int_{0}^{1} (1-\frac{z}{2}\cdot u)^{\tau-1}(1-u)^{\tau-1}du}$$
(3.14)

which is analytic for  $x \notin [-1,1]$ . Since the denominator  $\mathcal{D}_1(x)$  in (3.14) does not vanish in (-1,1)(it is  $X_0(x)$ , which does not vanish in (-1,1)), then

$$y_{1}(x) = \frac{1}{2\pi \lambda} \lim_{\varepsilon \to 0+} (X_{1}(x - \lambda \varepsilon) - X_{1}(x + \lambda \varepsilon))$$
 (3.15)

is continuous on R except, perhaps, at ±1. A simple but

rather long calculation gives for  $x \in (-1,1)$  that

$$y_{1}(x) = -\frac{1}{\pi} \cdot \text{Im } X_{1}(x)$$

$$= 2^{2\tau - 1} \cdot \frac{\Gamma(\tau)^{2}(\tau + 1)}{\pi \Gamma(2\tau) \tau} \cdot (1 - x^{2})^{\tau - \frac{1}{2}}$$
(3.16)

$$\frac{z[(z/\tau+1)F(1-\tau,2,\tau+2:z^2)-xF(1-\tau,1,\tau+1:z^2)]}{|F(1-\tau,1,\tau+1:z^2)|^2}$$

where

$$F(a,b,c:x) = \sum_{n=0}^{\infty} \frac{(a)_n(b)_n}{n!(c)_n} \cdot x^n.$$
 (3.17)

Other expressions for  $y_1(x)$  can be found in [4], [9], [22], [25]. We will write  $y_1(x) = (1-x^2)^{\tau-\frac{1}{2}}\tilde{y}(x)$ , where  $\tilde{y}(x) = y_1(x)(1-x^2)^{\frac{1}{2}-\tau}$  for  $x \in (-1,1)$  and  $\tilde{y}(x) = 0$  for  $x \notin [-1,1]$ . Then  $\tilde{Y}(x)$  is continuous for  $x \neq \pm 1$  and  $\tilde{Y}(x) \neq 0$  a.e. in (-1,1). The support Supp v of the orthogonality measure v of  $\{C_n^T(x;1)\}\$ is then [-1,1].

More details about the associated ultraspherical polynomials can be found in the references above.

§4. Basic relationships. Let  $\{P_n(x)\}$  be given by (1.1), (1.2) and (1.3). Eliminating  $P_{2n+1}(x)$  and  $P_{2n-1}(x)$  form (1.1) gives

$$(x^{2}-\frac{1}{2})P_{2n}(x) = P_{2(n+1)}(x) + \frac{1}{4^{2}} \cdot \frac{n}{n+\tau} \cdot \frac{n+2\tau-1}{n+\tau-1} P_{2(n-1)}(x)$$
 (4.1)

which hols for  $n \ge 2$ . Let

$$w = w(x) = 2x^2 - 1$$
 (4.2)

$$Q_{n}(x) = \frac{4^{n}(\tau)_{n}}{n!} \cdot P_{2n}(x), \quad n \ge 0.$$
 (4.3)

Then

Then
$$2(n+\tau)\omega Q_{n}(x) = (n+1)Q_{n+1}(x) + (n+2\tau-1)Q_{n-1}(x), \quad n \geqslant 2, \quad (4.4)$$

and from (2.9), (2.10) it follows that relupled good render

$$Q_{n}(x) = \alpha C_{n}^{\tau}(w) + 4\tau (1-\alpha) x^{2} C_{n-1}^{\tau}(w;1), \quad n \geqslant 1.$$
 (4.5)

Hence,

$$P_{2n}(x) = \frac{n!}{4^n(\tau)_n} \left[ \alpha C_n^{\tau}(W) + 4\tau (1-\alpha) x^2 C_{n-1}^{\tau}(W;1) \right], \quad n \ge 1.$$
 (4.6)

In particular (3.10) and (3.11) yield

$$P_{2n}(0) = (-1)^n \frac{(2\tau)_n}{4^n(\tau)_n} \cdot \alpha$$
 (4.7)

$$P_{2n}(\pm 1) = \frac{(2\tau)_n}{4^n(\tau)_n} \left[ \frac{4\tau - (2\tau + 1)\alpha}{2\tau - 1} \right] - \frac{4\tau(1 - \alpha)n!}{(2\tau - 1)4^n(\tau)_n}, \quad n > 1, \tau \neq \frac{1}{2}, \tag{4.8}$$

and

$$P_{2n}(\pm 1) = \frac{n!}{4^n (\frac{1}{2})_n} \left[ \alpha + 2(1-\alpha) \sum_{k=0}^{n-1} \frac{1}{k+1} \right], \quad n \ge 1, \tau = \frac{1}{2}.$$
 (4.9)

Let  $\mu$  denote the orthogonality measure of  $\{P_n(x)\}$ .

Then

Supp
$$\mu \subseteq \left[-(2+\sqrt{\alpha})/\sqrt{2},(2+\sqrt{\alpha})/\sqrt{2}\right]$$
 (4.10)

and

$$\int_{0}^{+\infty} P_{n}(x) P_{m}(x) d\mu = k_{n} \delta_{mn}; m, \quad m, n \geq 0, \quad (4.11)$$

with

$$k_{0} = 1, k_{1} = \frac{1}{2}$$

$$k_{2n} = \frac{\tau\alpha(2\tau)_{n} \cdot n!}{(n+\tau) \cdot 4^{2n} \cdot (\tau)_{n}^{2}}, k_{2n+1} = (\tau\alpha) \frac{(2\tau)_{n+1} \cdot n!}{(\tau)_{n+1}^{2} \cdot 4^{2n+1}}, \quad n \ge 1.$$
(4.12)

A calculation entirely similar to the one above gives for  $\{P_n^{(1)}(x)\}$  that

$$P_{2n+1}^{(1)}(x) = \frac{(n+1)!}{4^n(\tau+1)_n} x C_n^{\tau}(w;1), \quad n \ge 0.$$
 (4.13)

The continued fraction X(x) of  $\{P_n(x)\}$  then is

$$X(x) = \lim_{n \to \infty} \frac{P_{2n+1}^{(1)}(x)}{P_{2n+2}(x)} = \frac{2xX_{0}(w)}{\alpha + 2(1-\alpha)x^{2}X_{0}(w)}$$

$$x \neq \left[-(2+\sqrt{\alpha})/\sqrt{2}, (2+\sqrt{\alpha})/\sqrt{2}\right], \tag{4.14}$$

where  $X_0(W)$ , given by (3.12), is the continued fraction of  $\{C_n^T(W)\}$ . Hence X(x) is an analytic function of x for  $x \notin [-1,1]$ , except, perhaps, for simple poles located in  $[-(2+\sqrt{\alpha})/\sqrt{2},-1)$   $U(1,(2+\sqrt{\alpha})/\sqrt{2}]$ . Let

$$g(x) = \lim_{x \to 0} \frac{1}{2\pi i} \cdot \{X(x - i\varepsilon) - X(x + i\varepsilon)\}$$
 (4.15)

Then g(x) = 0 for  $x \notin [-(2+\sqrt{\alpha})/\sqrt{2}, (2+\sqrt{\alpha})/\sqrt{2}]$  and

$$g(x) = -\frac{2\alpha |x| \cdot \text{Im} X_0(w)}{\pi |\alpha + 2(1 - \alpha) x^2 X_0(w)|^2}, x \in (-1, 0) u(0, 1). \tag{4.16}$$

Let  $\mathcal{D}(x) = \alpha + 2(1-\alpha)x^2X_0(w)$ . Since

$$\operatorname{Im} X_{O}(W) = -2^{2\tau} \tau \frac{\Gamma(\tau)^{2}}{\Gamma(2\tau)} (1-W^{2})^{\tau-\frac{1}{2}}, \quad x \in (-1,0) U(0,1),$$
 (4.17)

as follows from (3.13), and

$$\operatorname{Im} \mathcal{D}(x) = 2(1-\alpha) x^{2} \cdot \operatorname{Im} X_{O}(w), \qquad (4.18)$$

 $\mathcal{D}(x)$  does not vanish on  $(-1,0)\,U(0,1)$ . An argument based on Remark 2.1 then shows that

$$\int_{-\infty}^{+\infty} \int d\mu = \int_{-\infty}^{+\infty} \int (x) g(x) dx \qquad (4.19)$$

for any continuous function with compact support contained in  $(-1;0)\,U(0,1)$ , that g is integrable on this set, and that this set is free of masses of  $\mu$ . Since g(x)>0 on  $(-1,0)\,U(0,1)$ , this set is part of  $C_\mu$ . Now,  $\mathcal{D}_\mu$  is the subset of  $\left[-(2+\sqrt{\alpha})/\sqrt{2},-1\right)\,U(1,(2+\sqrt{\alpha})/\sqrt{2}\right]$  where  $\mathcal{D}(x)$  vanishes, and  $P_\mu\subseteq\mathcal{D}_\mu U\{0,\pm 1\}$ . In the next section we will show that  $0\not=P_\mu$ , so that  $(-1,1)\subseteq C_\mu$ . We will also prove that  $\pm 1\subseteq P_\mu$  if and only if  $\tau>3/2$  and  $\alpha=4\tau/(2\tau+1)$ . Hence, if  $\alpha=1,C_\mu=[-1,1]$  and  $P_\mu=\phi$ , which agrees with results in [2].

Now we turn to  $\{P_n^{(1)}(x)\}$ . Since

$$xP_{2n}^{(1)}(x) = P_{2n+1}^{(1)}(x) + \frac{n+2\tau}{4(n+\tau)}P_{2n-1}^{(1)}(x), \quad n \geqslant 0,$$
 (4.20)

$$P_{2n}^{(1)}(x) = \frac{(n+1)!}{(\tau+1)_n 4^n} \left\{ C_n^{\tau}(w;1) + \frac{n+2\tau}{n+1} C_{n-1}^{\tau}(w;1) \right\}, \quad n \ge 0.$$
 (4.21)

From (2.10) and (2.11) we then obtain that

$$P_{2n}^{(1)}(0) = (-1)^n \frac{n!}{(\tau+1)_n 4^n}, \quad n \geqslant 0,$$
 (4.22)

$$P_{2n}^{(1)}(\pm 1) = \frac{2(2\tau + n)n!}{(2\tau - 1)(\tau + 1)_{n}4^{n}} \{ \frac{(2\tau)_{n}}{n!} - \frac{2n + 2\tau + 1}{4\tau + 2n} \}, \quad n \ge 0, \tau \pm \frac{1}{2}$$
 (4.23)

and

$$P_{2n}^{(1)}(\pm 1) = \frac{n!}{4^n(\tau+1)_n} \left\{ 2 \sum_{k=0}^{n-1} \frac{n+1}{k+1} + 1 \right\}, \quad n \ge 0, \tau = \frac{1}{2}. \tag{4.24}$$

The orthogonality relation for  $\{P_n^{(1)}(x)\}$  is

$$\int_{-\infty}^{+\infty} P_n^{(1)}(x) \cdot P_m^{(1)}(x) dv = k_n' \delta_{mn}; \ m, n \ge 0,$$
 (4.25)

where

$$k_{2n}' = \frac{(2\tau+1)_n n!}{(\tau+1)_n^2 4^{2n}}, \quad k_{2n+1}'' = \frac{(2\tau+1)_n (n+1)!}{(\tau+n+1)(\tau+1)_n^2 4^{2n+1}}, \qquad n \ge 0.$$
 (4.26)

For  $P_{2n}^{(2)}(x)$  we obtain, as before, that

$$P_{2n}^{(2)}(x) = \frac{(n+1)!}{4^n(\tau+1)_n} \left[ C_n^{\tau}(w;1) + \frac{1}{2} C_{n-1}^{\tau}(w;2) \right], \quad n \geqslant 0, \tag{4.27}$$

and the continued fraction for  $\{P_n^{(1)}(x)\}$  is

$$X^{*}(x) = \lim_{n \to \infty} \frac{p_{2n}^{(2)}(x)}{p_{2n+1}^{(1)}(x)} = \frac{1}{x} [1 + \frac{1}{2}X_{1}(w)], \qquad (4.28)$$

where  $X_1(w)$ , given by (4.13), is the continued fraction of  $\{C_n^{\tau}(W;1)\}$ . Hence  $X^*(x)$  is an analytic function of x for  $x \notin [-1.1]$ . We have that [-1,1]. We have that and P. = o, which agrees with rosust

$$y(x) = \lim_{\varepsilon \to 0+} \frac{1}{2\pi \lambda} \{ X^*(x - i\varepsilon) - X^*(x + i\varepsilon) \}$$

$$= \frac{2^{2\tau - 2}}{\pi} |x|^{2\tau - 2} (1 - x^2)^{\tau - \frac{1}{2}} \widetilde{y}(2x^2 - 1)$$
(4.29)

where  $\tilde{y}$  is given at the end of Section 3. As before we con-

clude that Y(x) is integrable on (-1,0) U (0,1), and it is continuous in (-1,0) U (0,1) if  $\tau > \frac{1}{2}$ . The absolutely continuous part of v is Y(x)dx, (-1,0) U  $(0,1) \subseteq C_v$ , and  $P_v \subseteq \{0,\pm 1\}$ . In the next section we will prove that  $P_v = \emptyset$  if  $\tau < \frac{1}{2}$  and  $P_v = \{0\}$  if  $\tau > \frac{1}{2}$ , so that  $C_v = [-1,1]$  or  $C_v = [-1,0)$  U (0,1], accordingly.

## §5. Main results. avail of seasons and an all and seasons

THEOREM 5.1. The orthogonality measure  $\mu$  of  $\{P_n(x)\}$  bears masses at  $x=\pm 1$  if and only if  $\tau>3/2$  and

$$\alpha = \frac{4\tau}{2\tau + 1} \,. \tag{5.1}$$

In such case, ±1 are end points masses of  $\mu$ .

**Proof.** From (4.8) and (4.12) we obtain, when  $\tau \neq \frac{1}{2}$ ,

that

$$\frac{P_{2n}(\pm 1)}{\sqrt{k_{2n}}} = \left[\frac{4\tau - (2\tau + 1)\alpha}{2\tau - 1}\right] \sqrt{\left[\frac{(n+\tau)(2\tau)_n}{n!\tau\alpha}\right]} + \left[\frac{4\tau(1-\alpha)}{2\tau - 1}\right] \sqrt{\left[\frac{n!(n+\tau)}{(2\tau)_n\tau\alpha}\right]}, \quad n \geqslant 1,$$

and we have, from (3.3), that

$$\sqrt{\left[\frac{(n+\tau)(2\tau)_n}{n!\tau\alpha}\right]} \approx \frac{n^{\tau}}{\sqrt{\left[\tau\alpha\Gamma(2\tau)\right]}}$$

and

$$\sqrt{\left[\frac{n!(n+\tau)}{(2\tau)_n\tau\alpha}\right]^2} \approx \sqrt{\left[\frac{\Gamma(2\tau)}{\tau\alpha}\right]n^{1-\tau}}.$$

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$$\frac{P_{2n}(\pm 1)}{\sqrt{k_{2n}}} \approx \frac{n^{\tau}}{\sqrt{[\tau \alpha \Gamma(2\tau)]}}$$

if  $\alpha = 1$  (even if  $\tau = \frac{1}{2}$ ), and also

$$\frac{P_{2n}(\pm 1)}{\sqrt{k_{2n}}} \approx \frac{2\tau - (2\tau + 1)\alpha}{2\tau - 1} \frac{n^{\tau}}{\sqrt{[\tau \alpha \Gamma(2\tau)]}}$$

if  $\alpha \neq 1$ ,  $\alpha \neq 4\tau/2\tau+1$  and  $\tau > \frac{1}{2}$ . In both cases  $\sum_{n=0}^{\infty} \frac{p_{2n}^2(\pm 1)}{k_{2n}}$ diverges and  $\mu$  has no mass points at ±1. Now assume  $\alpha = 4\tau/2\tau + 1.$ 

Then

$$\frac{P_{2n}(\pm 1)}{k_{2n}} \approx \sqrt{\left[\frac{2\alpha\Gamma(2\tau+2)}{\tau}\right]} n^{1-\tau},$$

even if  $\tau = \frac{1}{2}$ . If  $\tau < 3/2$ ,  $\sum_{n=0}^{\infty} P_{2n}^2(\pm 1)/k_{2n}$  is still divergent. But if  $\tau$  > 3/2, it is convergent, and  $\mu$  may have masses at ±1. To see that such is the case we have to show that  $\sum_{n=0}^{\infty} \frac{p_{2n+1}^2(\pm 1)}{k_{2n+1}}$  is convergent, but this follows at once from

$$\pm \frac{P_{2n+1}(\pm 1)}{k_{2n+1}} = \sqrt{\left[\frac{k_{2n+2}}{k_{2n+1}}\right]} \cdot \frac{P_{2n+2}(\pm 1)}{\sqrt{k_{2n+2}}} + \frac{n+2\tau}{4(n+\tau)} \sqrt{\left[\frac{k_{2n}}{k_{2n-1}}\right]} \frac{P_{2n}(\pm 1)}{\sqrt{k_{2n}}}$$

$$= \sqrt{2} \sqrt{\left[\frac{n+1}{n+\tau+1}\right]} \cdot \frac{P_{2n+2}(\pm 1)}{\sqrt{k_{2n+2}}} + \sqrt{2} \sqrt{\left[\frac{n+2\tau}{n+\tau}\right]} \cdot \frac{P_{2n}(\pm 1)}{\sqrt{k_{2n}}}.$$

All that is left to prove is that when  $\tau = \frac{1}{2}$  and  $\alpha = 4\tau/2\tau+1$ ,  $\sum_{n=0}^{\infty} \frac{P_{2n}^2(\pm 1)}{k_{2n}}$  diverges. But, in such case,

$$\frac{P_{2n}(\pm 1)}{\sqrt{k_{2n}}} = \sqrt{\left[\frac{2n+1}{\alpha}\right]} \{\alpha + 2(1-\alpha) \sum_{k=0}^{n-1} \frac{1}{k+1} \} \sim \sqrt{\left[2/\alpha\right]} \cdot 2(1-\alpha) n^{\frac{1}{2}} \log n$$

as follows from (4.9) and from 
$$\sum_{k=0}^{n-1} \frac{1}{k+1} \sim \log n, \quad n \to \infty,$$
 (5.2)

which is a consequence of the formula

$$\lim_{n \to \infty} \left\{ \sum_{k=0}^{n} \frac{1}{k+1} - \log n \right\} = c \tag{5.3}$$

for the Euler constant c (See [22], p.8.). Then  $\sum_{n=0}^{\infty} \frac{p_{2n}^2(\pm 1)}{k_{2n}}$ is divergent, and the proof is complete.

**THEOREM 5.2.** The orthogonality measure  $\mu$  of  $\{P_{\mu}(x)\}$ bears no mass at x = 0.

**Proof.** From (4.7) and (4.12)

$$\frac{P_{2n}(0)}{\sqrt{k_{2n}}} = \sqrt{\left[\frac{(2\tau)_n(n+\tau)\alpha}{\tau_n!}\right]} \approx \sqrt{\left[\frac{\alpha}{\tau\Gamma(2\tau)}\right]}n^{\tau}.$$

Since  $\tau > 0$ ,  $\sum_{n=0}^{\infty} \frac{p_{2n}^2(0)}{k_{2n}}$  diverges.

**THEOREM 5.3.** If  $\tau > \frac{1}{2}$ , and only in such case, the orthogonality measure  $\nu$  of  $\{P_n^{(1)}(x)\}$  carries a mass at x=0, which is then an embedded mass point of  $\nu$ .

**Proof.** In fact from (4.22) and (4.26),

$$\frac{p_{2n}^{(1)}(0)}{\sqrt{k_{2n}^{r}}} = (-1)^{n} \sqrt{\left[\frac{n!}{(2\tau+1)_{n}}\right]} \approx (-1)^{n} \sqrt{\left[\Gamma(2\tau)\right]} n^{-\tau}.$$

Hence,  $\sum_{n=0}^{\infty} \left[ \frac{P_{2n}^{(1)}(0)}{\sqrt{k_{2n}}} \right]^2 \text{ converges if and only if } \tau > \frac{1}{2}. \text{ Since } P_{2n+1}^{(1)}(0) = 0, \text{ this also holds for } \sum_{n=0}^{\infty} \left[ \frac{P_{n}^{(1)}(0)}{\sqrt{k_{n}}} \right]^2. \text{ Since } (-1,0) U(0,1) \subseteq C_v, 0 \text{ is in such case an embedded mass point } 06 \text{ v.} \quad \triangle$ 

**THEOREM 5.4.** The measure v bears no masses at  $x = \pm 1$ .

**Proof.** From (4.23) and (4.26) it follows that, for

$$\frac{P_{2n}^{(1)}(\pm 1)}{\sqrt{R_{2n}}} = \frac{2\tau}{2\tau - 1} \sqrt{\left[\frac{(2\tau + 1)_n}{n!}\right]} - \frac{2n + 2\tau - 1}{2\tau - 1} \sqrt{\left[\frac{n!}{(2\tau + 1)_n}\right]}.$$

But

$$\sqrt{\left[\frac{(2\tau+1)_n}{n!}\right]} \approx \frac{n^{\tau}}{\sqrt{\Gamma(2\tau)}}$$

and

$$\frac{2n+2\tau-1}{2\tau-1}\sqrt{\left[\frac{n!}{(2\delta-1)_n}\right]}\approx \frac{2}{2\tau-1}\sqrt{\left[\Gamma(2\tau)\right]}n^{1-\tau}.$$

Hence

$$\frac{P_{2n}^{(1)}(\pm 1)}{\sqrt{k_{2n}}} \approx \frac{2}{2\tau - 1} \cdot \frac{n^{\tau}}{\sqrt{\Gamma(2\tau)}}, \quad \tau > \frac{1}{2};$$

$$\frac{P_{2n}^{(1)}(\pm 1)}{\sqrt{k_{2n}^{2}}} \approx \frac{2\sqrt{\Gamma(2\tau)}}{2\tau - 1} \cdot n^{1 - \tau}, \quad \tau < \frac{1}{2}.$$

The measure v

In both cases  $\sum_{n=0}^{\infty} \left| \frac{p(1)(\pm 1)}{\sqrt{k^2 2n}} \right|^2$  is divergent. For  $\tau = \frac{1}{2}$  we get, from (4.24) and (5.2), that

$$\frac{P_{2n}^{(1)}(\pm 1)}{\sqrt{k_{2n}^{\prime}}} \approx n^{\frac{1}{2}} \log n,$$

$$\frac{\frac{P_{2n}^{(1)}(\pm 1)}{\sqrt{k_{2n}^{\prime}}} \approx n^{\frac{1}{2}} \log n,$$
and 
$$\sum_{n=0}^{\infty} \left| \frac{\frac{P_{2n}^{(1)}(\pm 1)}{\sqrt{k_{2n}^{\prime}}}}{\sqrt{k_{2n}^{\prime}}} \right|^2 \text{ also diverges.} \quad \blacktriangle$$

It follows, using (2.24), that

**THEOREM** 5.5. For  $\tau > \frac{1}{2}$  the measure v is

$$dv(x) = Y(x) dx + a\delta(x) dx$$
 (5.4)

where Y(x), given by (4.29), is integrable and  $\delta(x) dx$  is the Dirac measure at x = 0. Furthermore

$$a = v(\{0\}) = \frac{2\tau - 1}{2\tau}.$$
 (5.5)

REMARK 5.1. Observe that the measure v tends to concentrate on x = 0 when  $\tau \to +\infty$ . THEOREM S. 4. The measure a bears we make es

**REMARK 5.2.** When  $\tau = 1$ , (5.4) takes the particulary simple form

$$dv(x) = \frac{1}{\pi} \sqrt{[1-x^2]} \Phi(x) dx + \frac{1}{2} \delta(x) dx.$$
 (5.6)

where  $\Phi(x)$  is the characteristic function of [-1,1].

For  $\nu$ , the discrete support  $\mathcal{D}_{\nu}$  is empty. We now describe Du.

THEOREM 5.6. The discrete support Du of the orthogonality measure  $\mu$  of  $\{P_{\mu}(x)\}$  is a follows:

- (1) If  $\alpha \leq 1$ ,  $\mathcal{D}_{\mu} = \phi$ . The first  $\mathcal{D}_{\mu} = \phi$  is the first  $\mathcal{D}_{\mu} = \phi$  is the first  $\mathcal{D}_{\mu} = \phi$ . (2) If  $\alpha > 1$  and  $\tau \leqslant \frac{1}{2}$ ,  $\mathcal{D}_{\mu} = \{-c,c\}$ , where  $1 < c \leqslant (\sqrt{\alpha}+2)/\sqrt{2}$ is such that D(c) = 0.
- (3) If  $\tau > \frac{1}{2}$  and  $1 < \alpha \le 4\tau/(2\tau + 1)$ ,  $D\mu = \phi$ .

(4) If  $\tau > \frac{1}{2}$  and  $\alpha > 4\tau/2\tau+1$ ,  $D_{\mu} = \{-c,c\}$ , where, as before,  $1 < c \le (\sqrt{\alpha}+2)/\sqrt{2}$  is such that  $\mathcal{D}(c) = 0$ .

**Proof.** Let  $\mathcal{D}(x) = \alpha + 2(1-\alpha)x^2X_0(\emptyset)$ . Clearly  $\mathcal{D}(x) > 0$  if  $\alpha < 1$ . Let  $\mathcal{D}(+\infty) = \lim_{X \to +\infty} \mathcal{D}(x)$ ,  $\mathcal{D}(1^+) = \lim_{X \to 1^+} \mathcal{D}(x)$ . Then  $\mathcal{D}(+\infty) = 1$ , and

$$\mathcal{D}(1+) = -\infty, \tau \leq \frac{1}{2}; \mathcal{D}(1^+) = \frac{4\tau - (2\tau + 1)\alpha}{2\tau - 1}, \quad \tau > \frac{1}{2}$$

Hence  $\mathcal{D}\mu$  =  $\phi$  if  $\alpha < 1$ . If  $\alpha > 1, \mathcal{D}(1+)$  and  $\mathcal{D}(+\infty)$  have opposite signs in cases (2) and (4), and there must be  $1 < c < +\infty$  such that  $\mathcal{D}(c)$  = 0. Since the numerator  $\mathcal{N}(x)$  in (4.14) is positive for x > 0,  $\mu$  has a mass at x = c whose value is

$$\mu(\{c\}) = \text{Res}(X,c) = \frac{N(c)}{p'(c)}.$$
 (5.7)

Now, the set  $\mathcal{D}_{\mu}$  is discrete. If there were c' > c sich that  $(c,c') \cap \mathcal{D}_{\mu} = \phi$  and  $\mathcal{D}(c') = 0$  then  $\mathcal{D}'(c)$  and  $\mathcal{D}'(c')$  would carry opposite signs, which is absurd, since  $\mu(\{c'\}) > 0$ . Hence  $\mathcal{D}_{\mu} = \{-c,c\}$  in cases (2) and (4). Since it is impossible to have  $\mathcal{N}(x) = \mathcal{D}(x) = 0$  for x > 0, it follows that  $\mathcal{D}_{\mu} = \phi$  in case (3).

COROLLARY 5.1. If  $\tau > 3/2$  and  $\alpha = 4\tau/2\tau + 1$ , the orthogonality measure  $\mu$  of  $\{P_n(x)\}$  is

$$d\mu(x) = Y(x) dx + a [\delta(x-1) + \delta(x+1)] dx$$
 (5,8)

where Y(x), given by (4.16) for  $x \in (-1,0) \cup (0,1)$  and Y(x) = 0 elsewhere, is integrable,  $\delta(x\pm 1)$  is the Dirac measure at  $\pm 1$ , and

$$a = \frac{2\tau - 3}{4\tau - 2} . \tag{5.9}$$

**Proof.** Let  $\mathcal{D}(x) = \alpha + 2(1-\alpha)x^2X_0(\mathcal{W})$ ,  $x \in [-1,1]$ ,  $\mathcal{W} = 2x^2-1$ , and extend  $\mathcal{D}(x)$  to (-1,1) by taking  $X_0(\mathcal{W}(x)) = \lim_{\varepsilon \to 0+} X_0(\mathcal{W}(x+i\varepsilon))$  for  $x \in (-1,1)$ . Then  $\mathcal{D}(x)$  is anlytic in  $\mathbb{C}$ -[-1,1] and its restriction to (-1,1) is also continuous. It is readily seen that for  $\tau > 3/2$  and  $\alpha = 4\tau/2\tau+1$  we have

$$\mathcal{D}(x) = (1-x^2)F(x), \quad x \in [-1,1],$$

where F(x) is continuous and non vanishing in (-1,1), and

$$F(\pm 1) = -4\tau/2\tau - 3$$
, (x) I set aloom

so that

$$Y(x) = \begin{cases} 4^{2\tau+1} \cdot \frac{\Gamma(\tau)^2}{\Gamma(2\tau)} \cdot \frac{|x|^{2\tau} (1+x)^{\tau-5/2} (1-x)^{\tau-5/2}}{|F(x)|^2}, & x \in (-1,1) \\ 0, & x \notin [-1,1], \end{cases}$$

which obviously is integrable. (a) I said dollar

To show that a is given by (5.9), let  $\varepsilon, \delta > 0$  and integrate over the positively oriented rectangular path  $\beta_{\varepsilon,\delta}$  with center at 1, basis  $2\delta$  and height  $2\varepsilon$ . We have

$$\int_{1-\delta}^{1+\delta} \left[ X(x-i\varepsilon) - X(x+i\varepsilon) \right] dx + i \int_{-\varepsilon}^{\varepsilon} \left[ X(1+\delta+it) - X(1-\delta-it) \right] dt = \int_{\beta\varepsilon, \delta} X(x) dx$$

Now, the right hand side integral can be replaced by the integral on  $\beta_{\delta}$ , the positively oriented circle with center at 1 and radius  $\delta$  (because X(x) is analytic on Im x > 0 and Im x < 0 and both  $\lim_{\varepsilon \to 0} X(x \pm i\varepsilon)$  exist and are finite for  $x \in [1-\delta,1)U(1,1+\delta]$ ). Moreover, the second integral on the left vanishes when  $\varepsilon \to 0$ . Hence

$$\sigma(1+\delta)-\sigma(1-\delta) = \lim_{\varepsilon \to 0} \frac{1}{2\pi \lambda} \int_{1-\delta}^{1+\delta} (X(x-i\varepsilon)-X(x+i\varepsilon)) dx = \frac{1}{2\pi \lambda} \int_{\beta_{\delta}}^{1+\delta} X(x) dx,$$

and V White (I, The (in acres by x

$$a = \mu(\{1\}) = \sigma(1) - \sigma(1-0) = \lim_{\delta \to 0} \frac{1}{2\pi \lambda} \int_{\beta_{\delta}} X(x) dx = -\frac{4\tau/2\tau - 1}{2F(1)} = \frac{2\tau - 3}{2(2\tau - 1)}.$$

The proof is complete. A

**REMARK 5.3.** Observe that  $\mu$  tends to concentrate on  $\pm 1$  when  $\tau \to +\infty$ .

#### §6. Final. Remarks.

Let  $\{p_n(x)\}$  be an O.N.P.S. given by

$$xp_n(x) = b_n p_{n-1}(x) + a_n p_n(x) + b_{n+1} p_{n+1}(x), \quad n > 0,$$
 (6.1)  
 $p_{-1}(x) = 0, \quad p_0(x) = 1.$ 

Assume

$$|a_n| \le M, 0 < b_{n+1} \le M, \quad n \ge 0.$$
 (6.2)

(bo is arbitrary). The tridiagonal infinite matrix

is called the Jacobi Matrix of  $\{p_n(x)\}$ . It defines on the Hilbert space  $\ell_2(\mathfrak{C})$  of square summable, complex sequences  $\{x_n/n > 0\}$  a bounded self-adjoint opperator  $\ell$  by

$$L(e_n) = b_n e_{n-1} + a_n e_n + b_{n+1} e_{n+1}, \quad n \ge 0,$$
 (6.4)

 $(e_{-1} = 0)$  and continuous linear extension. Here  $\{e_n/n \ge 0\}$  is the canonical basis of  $\ell_2(\mathbf{C})$   $(e_n = (\delta_{0n}, \delta_{1n}, \ldots), n \ge 0)$ . It can be shown that  $\|L\| \le 3M$ . Observe that  $x = (x_0, x_1, \ldots)$  is an eigenvector of L for the eigenvalue  $\beta$  if and only if  $\sum_{k=0}^{\infty} |x_k|^2 < \infty$  and k=0

$$\beta x_n = b_n x_{n-1} + a_n x_n + b_{n+1} x_{n+1}, \quad n \ge 0, \tag{6.5}$$

i.e, if and only if  $x \in \ell_2(\mathfrak{C})$  and  $x_n = x_0 p_n(\beta)$ , n > 0.

Let  $\mu$  be the orthogonality measure of  $\{p_n(x)\}$  and for each  $\tau \in \mathbb{R}$  let  $E_n: \ell_2 \to \ell_2$  be defined by

$$E_{\tau} e_{n} = \sum_{k=0}^{\infty} \{ \int_{-\infty}^{\tau} p_{n}(t) p_{k}(t) d\mu(t) \} e_{k}, \qquad (6.6)$$

and continuous linear extension. Then  $\{E_{\tau} | \tau \in \mathbb{R}\}$  is a resolution of the identity for L and

$$\int_{-\infty}^{+\infty} f(x) d\mu = \int_{-\infty}^{+\infty} f(t) d(E_{\tau} e_{0}; e_{0})$$
(6.7)

for any continuous function on  $\mathbb{R}((x;y) = \sum_{n=0}^{\infty} x_n \bar{y}_n$  is the inner product in  $\ell_2(\mathfrak{C})$ ). From (6.7) it easily follows that the spectrum  $\sigma(L)$  of L coincides with Supp  $\mu;\sigma_{\mathfrak{C}}(L)$ , the continuous spectrum, is  $C_{\mu}$ ; and  $\sigma_{\mathfrak{p}}(L)$ , the point spectrum, is  $P_{\mu}$ . Details of above can be found in [1], [6].

It follows that the operator L of the Jacobi matrix

$$a_n = 0$$
,  $n \ge 0$   
 $b_1 = \sqrt{[\alpha/2]}$ ,  $b_{2n} = \frac{1}{2}\sqrt{\frac{n}{n+\tau}}$ ,  $b_{2n+1} = \frac{1}{2}\sqrt{\frac{n+2\tau}{n+\tau}}$ ,  $n \ge 1$ 

has, for  $\tau > 3/2$  and  $\alpha = 4\tau/2\tau + 1$ , eigenvalues at the end points  $\pm 1$  of the spectrum  $\sigma(L) = [-1,1]$ . These are the only eigenvalues of L. As for the operator L of the Jacobi matrix

$$a_n = 0$$
,  $n \ge 0$  (6.9)  
 $b_{2n} = \frac{1}{2}\sqrt{\left[\frac{n+2\tau}{n+\tau}\right]}$ ,  $b_{2n+1} = \frac{1}{2}\sqrt{\left[\frac{n+1}{n+\tau+1}\right]}$ ,  $n \ge 0$ ,

it has, for  $\tau > \frac{1}{2}$ , an embedded eigenvalue, 0, in the spectrum  $\sigma(L) = [-1,1]$ . This is its only eigenvalue.

Now consider the recurrence relation

$$xp_{kn+j}(x) = p_{kn+j+1}(x) + a_{j,n}p_{nk+j-1}(x), n \ge 0, \quad j = 0,...,k-1,$$

$$p_{-1}(x) = 0, \quad p_{0}(x) = 1,$$
(6.10)

where  $k \ge 2$  is an integer,  $a_{0,n}$  and  $a_{1,n}$ ,  $n \ge 0$ , are given by (1.3), and

$$a_{j,n} = 1/4, j = 2,3,...,k-1, n > 0.$$
 (6.11)

Recurrence relation (6.10) reduces to that of the sieved

ultraspherical polynomials of the first kind in [2] when  $\alpha = 1$ . As it is to be expected, the orthogonality measure  $\mu$  of  $\{p_n(x)\}$  bears masses at  $\pm 1$  when  $k \ge 2$ ,  $\tau > 3/2$  and  $\alpha = 4\tau/2\tau+1$ . When  $\tau > \frac{1}{2}$ , the orthogonality measure of  $\{p_n^{(1)}(x)\}$  carries embedded masses at the points x of (-1,1) where  $T_k(x) = \pm 1(T_k(x))$  is the  $k^{th}$  Chebichev polynomial of the first kind ([22], [25])). The technical details of the proofs will appear elsewhere.

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